

Bachelor in Aerospace Engineering



Universidad
Carlos III de Madrid
www.uc3m.es

Department of Continuum Mechanics and Structural Engineering

Aerospace Structures

Chapter 4. Laminate and sandwich structures

Composite Beams



Chapter 4. Composite beams and plates

Composite beams

- 1. Introduction**
- 2. Analysis of beams**
- 3. Stresses in solid beams**
- 4. Displacements in solid beams**
- 5. Stresses in thin-walled beams**
- 6. Displacements in thin-walled beams**
- 7. References**

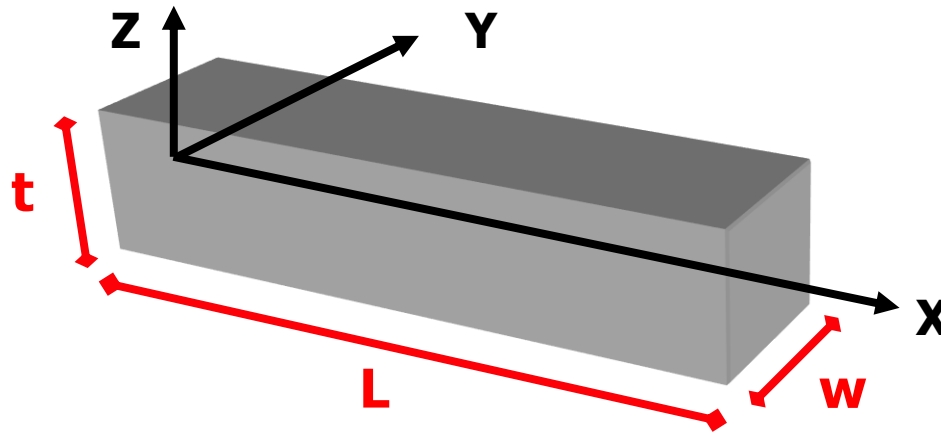


Chapter 4. Composite beams and plates

Composite beams

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Previous knowledge



L = Length

t = Thickness

w = Width

- $w / L \ll 1$
- $t / L \ll 1$
- Loads directions are in plane x-z.
- If loads directions are in plane y-z, we can apply analogous formulation.
- Principle of superposition can be use for loads direction in plane x-z and plane y-z.



Internal force in beams:

- Tensile
- Compressive
- Bending moment
- Torsion moment
- Shear
- Combined forces:
Bending-shear, bending-torsion, ...

Manufacture methods

- Pultrusion
- Rolling
- Prepegç
- Resin injection



- Design:
 - Minimize cross-section (weight)
 - Maximize moments of inertia
 - Increase section dimensions
 - Decrease thickness
 - Beams modelled as prismatic elements
- Design stages:
 - Beam dimensions, under quasi-static or dynamic
 - Joint of beam with the rest of the structure



Chapter 4. Composite beams and plates

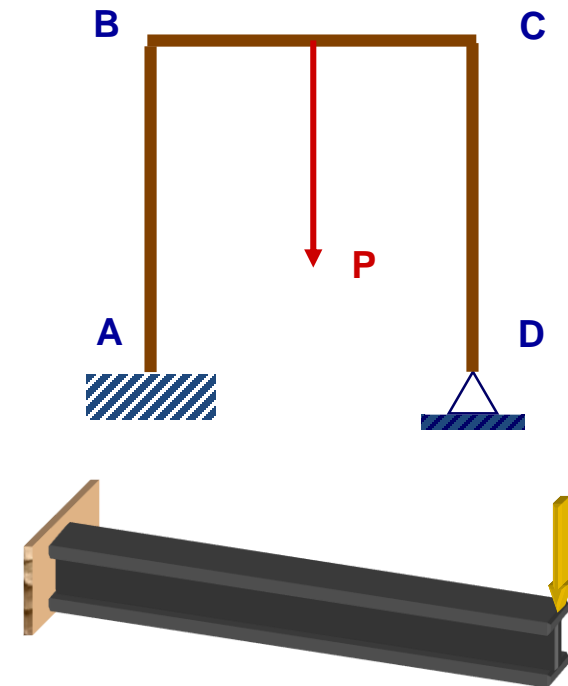
Composite beams

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Calculation process

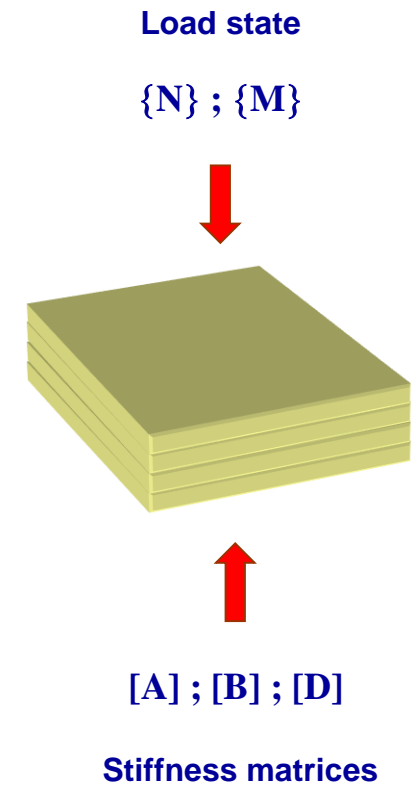
– Classical beam theory can be used to analyse the structure

- Degree of static indeterminacy
- Reactions calculation
- Internal forces laws



Calculation process

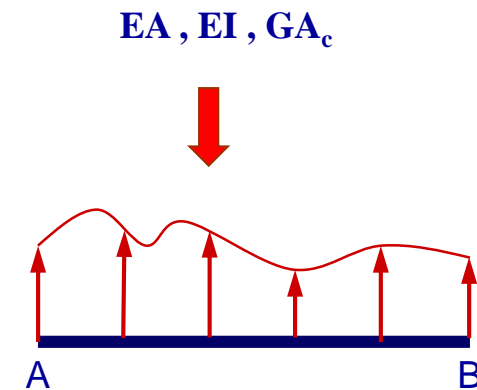
- Stresses analysis:
 - Laminate theory or first-order-shear-deformation theory can be used to calculate beam strains
 - Strains and stresses can be found for each layer
 - A failure criterion is applied



Calculation process

– Displacement analysis

- Laminate equivalent elastic properties are found using laminate theory or first-order-shear-deformation theory to calculate stiffness matrices
- Navier-Bresse formulation, Mohr theorems or energetic methods can be used to find the displacements



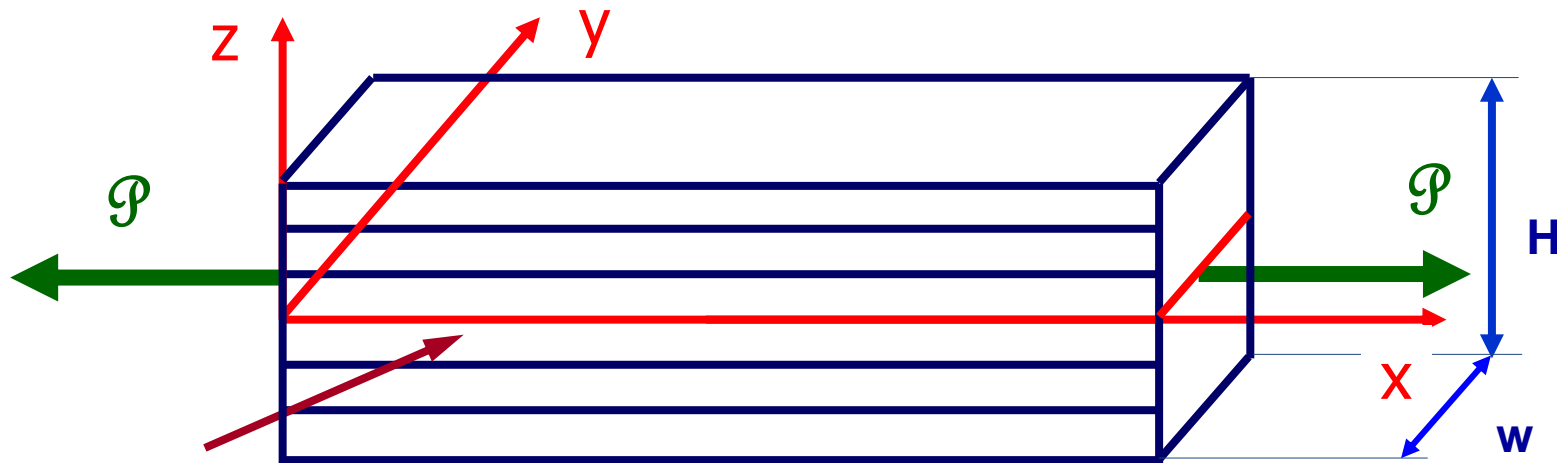


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Beams under tensile/compressive loads



Mid-plane

\mathcal{P} : External load (N)

N : Internal axial force per unit length (N/m)

$$N_x = N = \frac{\mathcal{P}}{w}$$

Beams under tensile/compressive loads

Laminate theory:

$$\{N\} = \underbrace{\left[\int_{-t/2}^{t/2} [\bar{Q}] \cdot dz \right]}_{[A]} \cdot \{\varepsilon^0\}$$

$$\{N\} = [A] \cdot \{\varepsilon^0\} \text{ in N/m}$$

Beams under tensile/compressive loads

Load state:

$$\begin{Bmatrix} N \\ 0 \\ 0 \end{Bmatrix} = \begin{bmatrix} A_{11} & A_{12} & A_{1S} \\ A_{12} & A_{22} & A_{2S} \\ A_{1S} & A_{2S} & A_{SS} \end{bmatrix} \begin{Bmatrix} \varepsilon_x^0 \\ \varepsilon_y^0 \\ \gamma_{xy}^0 \end{Bmatrix}$$

Strains:

$$\begin{Bmatrix} \varepsilon_x^0 \\ \varepsilon_y^0 \\ \gamma_{xy}^0 \end{Bmatrix} = \begin{bmatrix} a_{11} & a_{12} & a_{1S} \\ a_{12} & a_{22} & a_{2S} \\ a_{1S} & a_{2S} & a_{SS} \end{bmatrix} \begin{Bmatrix} N \\ 0 \\ 0 \end{Bmatrix}$$

$$\varepsilon_x^0 = a_{11} \cdot N$$

$$N = N(x)$$

$$\varepsilon_x^0 = \varepsilon_x^0(x)$$

One-dimensional problem

Beams under tensile/compressive loads

Strains of each lamina

Lamina k-th:

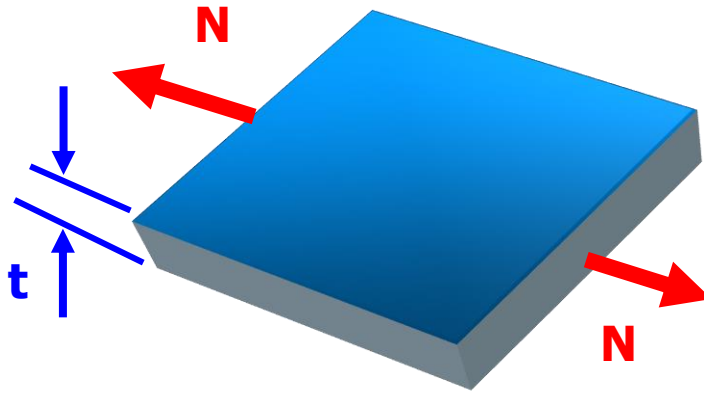
$$\{\boldsymbol{\varepsilon}\}^k = \begin{Bmatrix} \boldsymbol{\varepsilon}_x \\ \boldsymbol{\varepsilon}_y \\ \boldsymbol{\gamma}_{xy} \end{Bmatrix}^k = \begin{Bmatrix} \boldsymbol{\varepsilon}_x^0 \\ \boldsymbol{\varepsilon}_y^0 \\ \boldsymbol{\gamma}_{xy}^0 \end{Bmatrix} = \{\boldsymbol{\varepsilon}^0\}$$

Stresses in each lamina

$$\{\boldsymbol{\sigma}\}^k = [\bar{Q}]_k \{\boldsymbol{\varepsilon}\}^k = [\bar{Q}]_k \{\boldsymbol{\varepsilon}^0\}$$

Beams under tensile/compressive loads

Laminate equivalent constants



$$[a^*] = \frac{1}{[A^*]} \quad [A^*] = \frac{[A]}{t}$$

$$\varepsilon_x^0 = a_{11} \cdot N \quad \sigma_x^0 = \frac{N}{t} = \frac{\mathcal{P}}{t \cdot w}$$

In isotropic materials

$$\sigma_x = E \cdot \varepsilon_x \quad \sigma_x = \frac{\mathcal{P}}{A}$$



$$E_x^0 = \frac{1}{a_{11} \cdot t} = \frac{1}{a_{11}^*} \quad A = w \cdot t$$

$$\langle EA \rangle = \frac{w}{a_{11}}$$

Beams under tensile/compressive loads

Approximate solution

$$\{\varepsilon\}^k = \{\varepsilon^0\} = [\bar{S}]_k \{\sigma\}^k$$

$$\sigma_x^k = E_x^k \varepsilon_x^0$$

$$\varepsilon_x^0 = \frac{\sigma_x^k}{E_x^k}$$

$$\sigma_x^k \cdot t_k = E_x^k \cdot t_k \cdot \varepsilon_x^0$$

$$N = \sum_k \sigma_x^k \cdot t_k = \left(\sum_k E_x^k \cdot t_k \right) \cdot \varepsilon_x^0$$

Beams under tensile/compressive loads

Where:

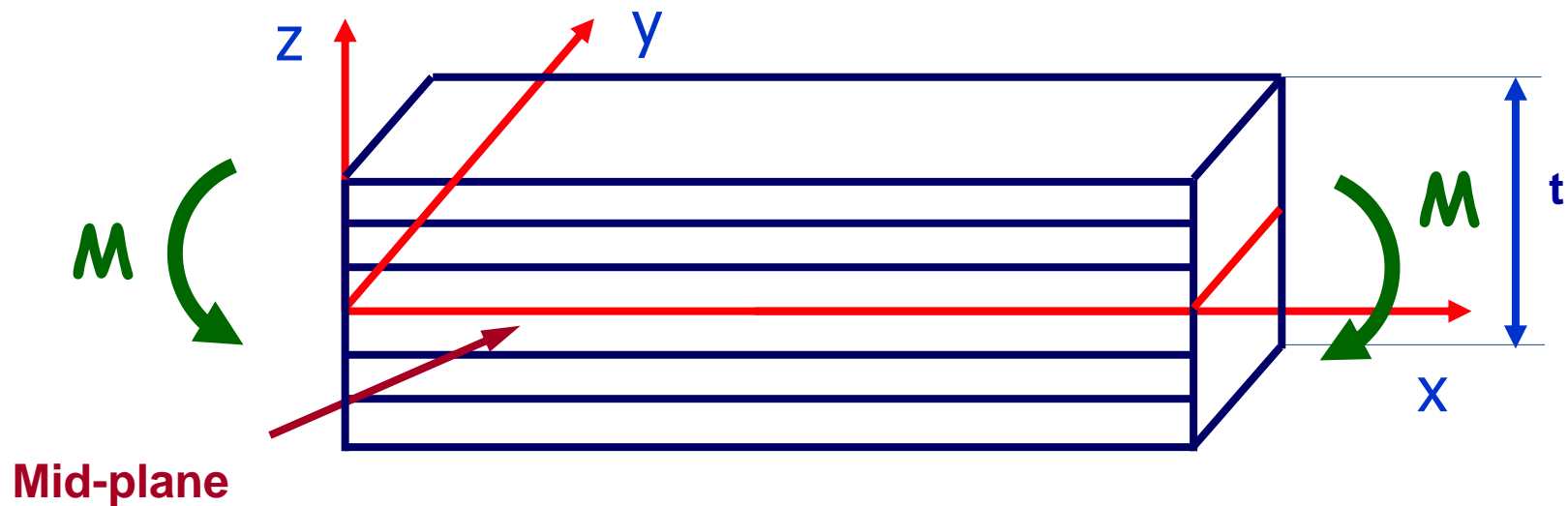
$$E_x^0 = \frac{\sum_k E_x^k \cdot t_k}{t}$$

$$N = \sum_k \sigma_x^k \cdot t_k = \left(\sum_k E_x^k \cdot t_k \right) \cdot \varepsilon_x^0 = E_x^0 \cdot t \cdot \varepsilon_x^0$$

$$\sigma_x^0 = \frac{N}{t} = E_x^0 \cdot \varepsilon_x^0$$

$$\varepsilon_x^0 = \frac{1}{E_x^0} \frac{N}{t}$$

Bending of beams



M : External moment (N·m)

$$M_x = M = \frac{M}{w}$$

M_x : Bending moment per unit length (N)

Bending of beams

(Narrow beams)

Load state:

$$\{M\} = \begin{Bmatrix} M_x \\ M_y \\ M_{xy} \end{Bmatrix} = \begin{Bmatrix} M \\ 0 \\ 0 \end{Bmatrix}$$

Laminate theory:

Symmetric laminate
under pure bending
moment

$$\begin{Bmatrix} M \\ 0 \\ 0 \end{Bmatrix} = \begin{bmatrix} D_{11} & D_{12} & D_{1S} \\ D_{12} & D_{22} & D_{2S} \\ D_{1S} & D_{2S} & D_{SS} \end{bmatrix} \begin{Bmatrix} \kappa_x \\ \kappa_y \\ \kappa_{xy} \end{Bmatrix}$$

Bending of beams

Laminate curvatures:

$$\begin{Bmatrix} \kappa_x \\ \kappa_y \\ \kappa_{xy} \end{Bmatrix} = \begin{bmatrix} d_{11} & d_{12} & d_{1S} \\ d_{12} & d_{22} & d_{2S} \\ d_{1S} & d_{2S} & d_{SS} \end{bmatrix} \begin{Bmatrix} M \\ 0 \\ 0 \end{Bmatrix}$$

$$K_x = d_{11} \cdot M$$

$$M = M(x)$$

$$k_x = k_x(x)$$

**One-dimensional
problem**

Bending of beams

Strains in each lamina

Lamina k-th:

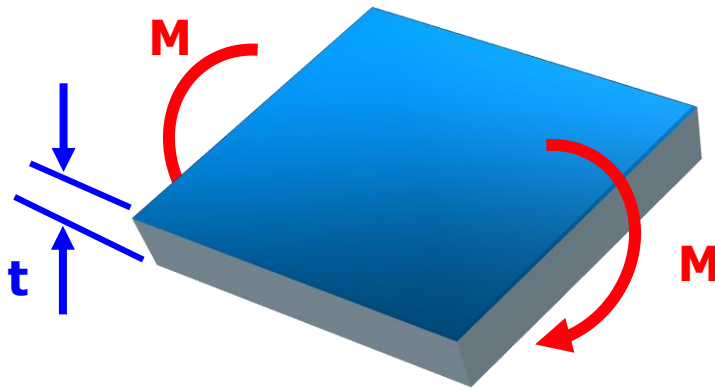
$$\{\varepsilon\}^k = \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix}^k = z \cdot \begin{Bmatrix} k_x \\ k_y \\ k_{xy} \end{Bmatrix} = z \cdot \{k\}$$

Stresses in each lamina

$$\{\sigma\}^k = [\bar{Q}]_k \{\varepsilon\}^k$$

Bending of beams

Laminate equivalent constants



$$[d^*] = \frac{1}{[D^*]} \quad [D^*] = \frac{12 \cdot [D]}{t^3}$$

$$\kappa_x = d_{11} \cdot M = \frac{1}{\rho_x}$$

In isotropic materials:

$$\frac{1}{\rho_x} = \frac{M}{E \cdot I}$$

$$I_y = \frac{1}{12} \cdot w \cdot t^3$$

$$E_x^b = \frac{12}{t^3 \cdot d_{11}} = \frac{1}{d_{11}^*}$$

$$\langle EI \rangle = \frac{w}{d_{11}}$$

Bending of beams

Laminate average maximum stresses

In symmetrical laminates under bending an equivalent maximum stresses can be defined as:

$$\sigma_x^b = \frac{6M}{t^2} \quad \sigma_y^b = 0 \quad \tau_x^b = 0$$


$$\{\sigma\}^b = \begin{Bmatrix} \frac{6M}{t^2} \\ 0 \\ 0 \end{Bmatrix} = \frac{6}{t^2} \{M\}$$

Bending of beams

Laminate maximum strains

Equivalent maximum strains:

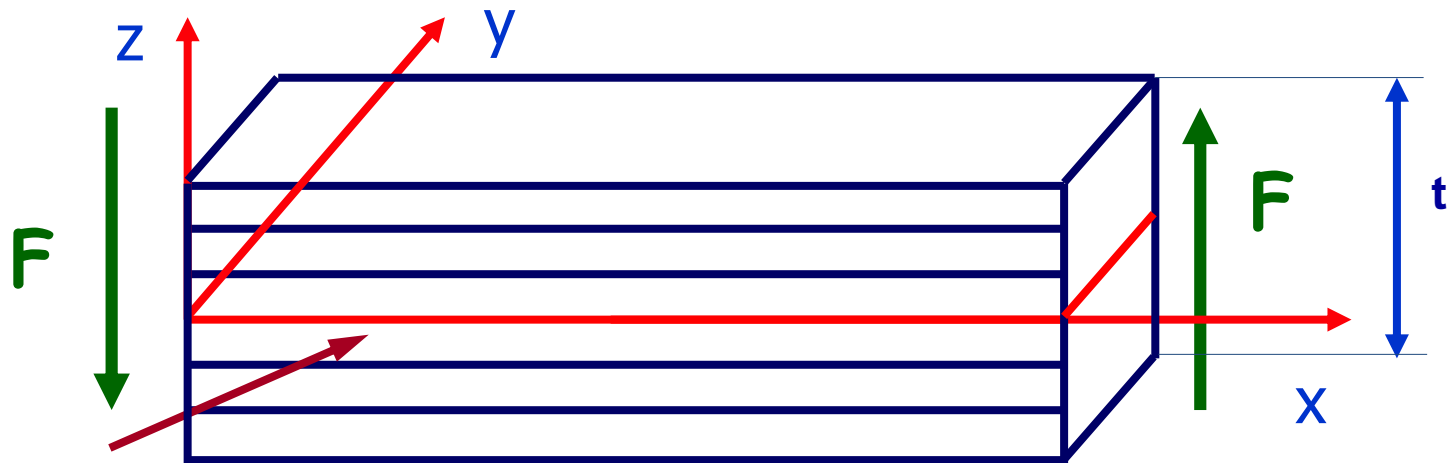
$$\{\varepsilon\}^b = \frac{t}{2} \cdot \{\kappa\}$$
$$\{\varepsilon\}^b = \frac{t}{2} \cdot \{\kappa\} = \frac{t}{2} [D]^{-1} \{M\}$$



$$\left\{ \begin{array}{l} \varepsilon_x^b = \frac{t}{2} d_{11} M \\ \varepsilon_y^b = \frac{t}{2} d_{12} M \\ \gamma_{xy}^b = \frac{t}{2} d_{1s} M \end{array} \right.$$

Since a linear through-the-thickness strains variation is assumed, strains at any layer can be found. Therefore stresses can be also found.

Interlaminar shear forces in beams



Mid-plane

F : External force (N)

N_{xy} : Shear force per unit length (N/m)

$$Q_x = Q = \frac{F}{w}$$

Interlaminar shear forces in beams

Load state:

$$\{Q\} = \begin{Bmatrix} Q_y \\ Q_x \end{Bmatrix} = \begin{Bmatrix} 0 \\ Q \end{Bmatrix}$$

First-Order Shear Deformation theory:

$$\begin{Bmatrix} 0 \\ Q \end{Bmatrix} = K \begin{Bmatrix} A_{44} & A_{45} \\ A_{54} & A_{55} \end{Bmatrix} \begin{Bmatrix} \gamma_{yz} \\ \gamma_{xz} \end{Bmatrix}$$

Strains:

$$\begin{Bmatrix} \gamma_{yz} \\ \gamma_{xz} \end{Bmatrix} = \frac{1}{K} \begin{Bmatrix} a_{44} & a_{45} \\ a_{54} & a_{55} \end{Bmatrix} \begin{Bmatrix} 0 \\ Q \end{Bmatrix} \quad \gamma_{xz} = \frac{a_{55}}{K} \cdot Q$$

**One-Dimensional
problem**

Interlaminar shear forces in beams

Lamina strains

Lamina k-th:

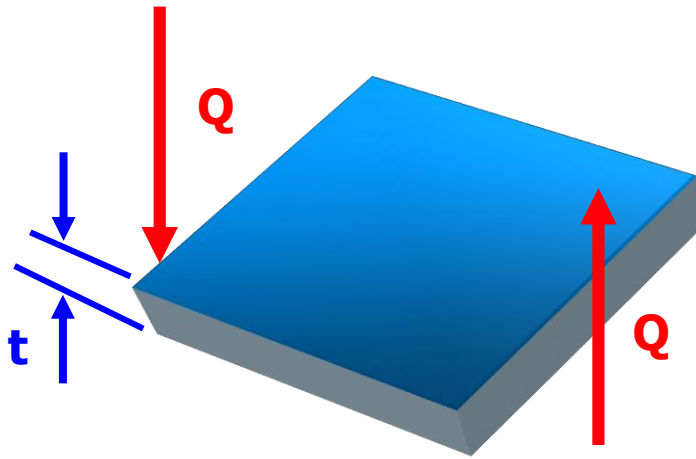
$$\{\gamma\}^k = \begin{Bmatrix} \gamma_{yz} \\ \gamma_{xz} \end{Bmatrix}^k = \begin{Bmatrix} \gamma_{yz}^0 \\ \gamma_{xz}^0 \end{Bmatrix} = \{\gamma^0\}$$

Lamina stress:

$$\{\tau\}^k = [\bar{Q}]_k \cdot \{\gamma\}^k = [\bar{Q}]_k \cdot \{\gamma^0\}$$

Interlaminar shear forces in beams

Laminate equivalent constants



$$\gamma_{xz}^0 = \frac{a_{55}}{K} \cdot Q \qquad \tau_{xz}^o = \frac{Q}{t}$$

In isotropic materials:

$$\tau_{xz} = G \cdot \gamma_{xz} \qquad \tau_{xz} = \frac{F}{A_s}$$

$$G_{xz}^0 = \frac{1}{a_{55} \cdot t} \qquad A_s = K \cdot w \cdot t$$

$$\langle GA_s \rangle = \frac{K \cdot w}{a_{55}}$$



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Displacements in solid beams

■ Long beams

- $L/H > 10$
- The shear force contribution to the displacements can be neglected. Displacements are mainly produced by bending moment

$$d = d_b$$

d_b controlled by EI

■ Short beams

- $L/H < 10$
- The shear force contribution to the displacements can not be neglected

$$d = d_b + d_s$$

d_b controlled by EI
 d_s controlled by GA

Calculation process

- Classical laminate theory is used to find equivalent stiffness to bending moments (E_b) and shear forces (G_{xz})
- Navier-Bresse formulation can be used to find the displacements

$$\theta_B = \theta_A - \int_A^B \frac{M}{EI} \cdot ds$$

$$v_B = v_A + \theta_A \cdot (z_B - z_A) + \int_A^B \left(\frac{N}{EA} dy + \frac{Q}{GA_s} dz \right) - \int_A^B \frac{M}{EI} (z_B - z) \cdot ds$$

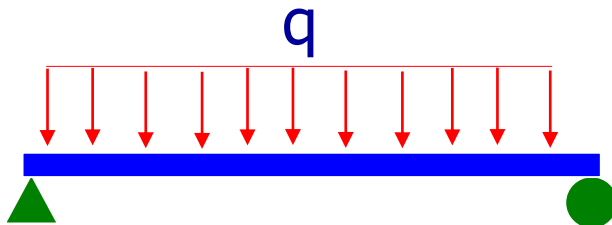
$$w_B = w_A + \theta_A \cdot (y_B - y_A) + \int_A^B \left(\frac{N}{EA} dz - \frac{Q}{GA_s} dy \right) + \int_A^B \frac{M}{EI} (y_B - y) \cdot ds$$

- Where

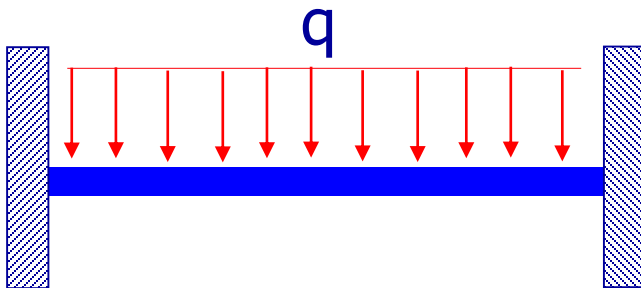
$$E \rightarrow E_b$$

$$G \rightarrow G_{xz}$$

Maximum displacements

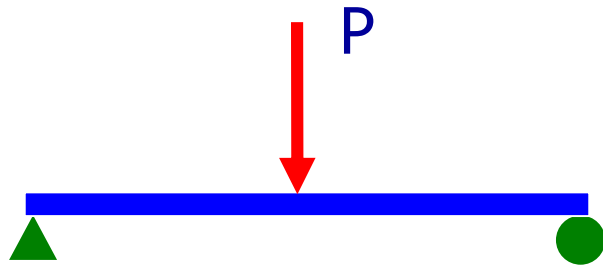


$$\delta_{\max} = \frac{5}{384} \frac{qL^4}{\langle EI \rangle} + \frac{1}{8} \frac{qL^2}{\langle GA_s \rangle}$$

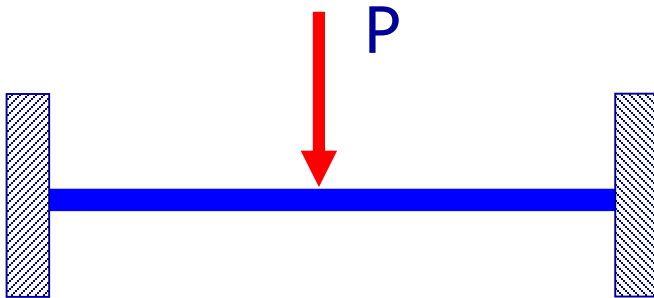


$$\delta_{\max} = \frac{1}{384} \frac{qL^4}{\langle EI \rangle} + \frac{1}{8} \frac{qL^2}{\langle GA_s \rangle}$$

Maximum displacements



$$\delta_{\max} = \frac{1}{48} \frac{PL^3}{\langle EI \rangle} + \frac{1}{4} \frac{PL}{\langle GA_s \rangle}$$



$$\delta_{\max} = \frac{1}{192} \frac{PL^3}{\langle EI \rangle} + \frac{1}{4} \frac{PL}{\langle GA_s \rangle}$$

Displacements in solid beams

Maximum displacements

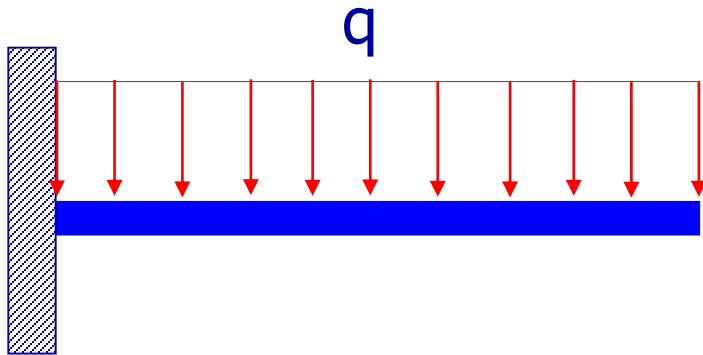


$$\delta_{\max} = \frac{ML^2}{8\langle EI \rangle}$$



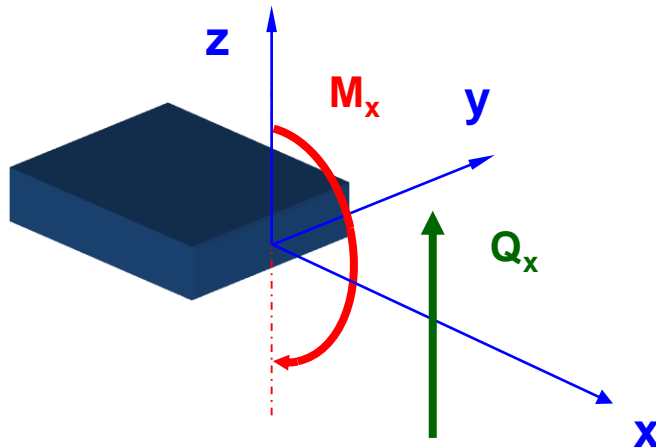
$$\delta_{\max} = \frac{PL^3}{3\langle EI \rangle} + \frac{\cancel{PL}}{\cancel{\langle GA_s \rangle}}$$

Maximum displacements



$$\delta_{\max} = \frac{1}{8} \frac{qL^4}{\langle EI \rangle} + \frac{1}{2} \frac{qL^2}{\langle GA_s \rangle}$$

Analysis of beams



$$M_y = 0$$

$$M_{xy} = 0$$

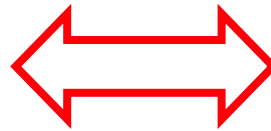
$$Q_y = 0$$

$$\phi_y = 0$$

Analysis of beams

$$\begin{Bmatrix} M_x \\ M_y \\ M_{xy} \end{Bmatrix} = \begin{bmatrix} D_{11} & D_{12} & D_{1s} \\ D_{12} & D_{22} & D_{2s} \\ D_{1s} & D_{2s} & D_{ss} \end{bmatrix} \cdot \begin{Bmatrix} K_x \\ K_y \\ K_{xy} \end{Bmatrix}$$

$$\{k\} = [d]\{M\}$$



$$k_x = \frac{\partial \phi_x}{\partial x}$$

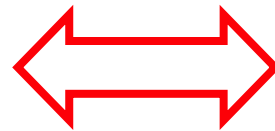
$$k_y = \frac{\partial \phi_y}{\partial y}$$

$$k_{xy} = \frac{\partial \phi_x}{\partial y} + \frac{\partial \phi_y}{\partial x}$$

Analysis of beams

$$\begin{Bmatrix} Q_y \\ Q_x \end{Bmatrix} = K \cdot \begin{Bmatrix} A_{44} & A_{45} \\ A_{54} & A_{55} \end{Bmatrix} \begin{Bmatrix} \gamma_{yz} \\ \gamma_{xz} \end{Bmatrix}$$

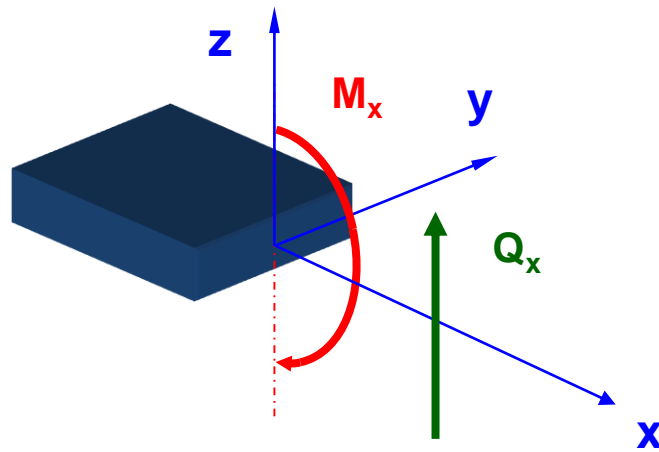
$$\{\gamma\} = \frac{1}{K} [a] \{Q\}$$



$$\gamma_{yz} = \frac{\partial w}{\partial y} + \phi_y$$

$$\gamma_{xz} = \frac{\partial w}{\partial x} + \phi_x$$

Analysis of beams



$$\frac{\partial \phi_x}{\partial x} = d_{11} M_x$$

$$\frac{\partial w}{\partial x} + \phi_x = \frac{a_{55}}{K} Q_x$$

$$M_y = 0 \quad Q_y = 0$$

$$M_{xy} = 0 \quad \phi_y = 0$$

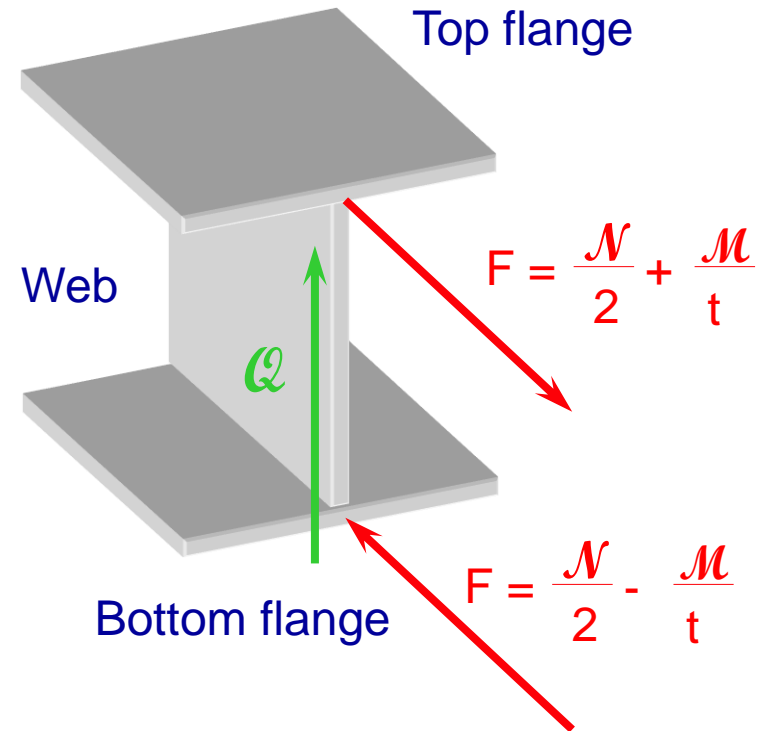
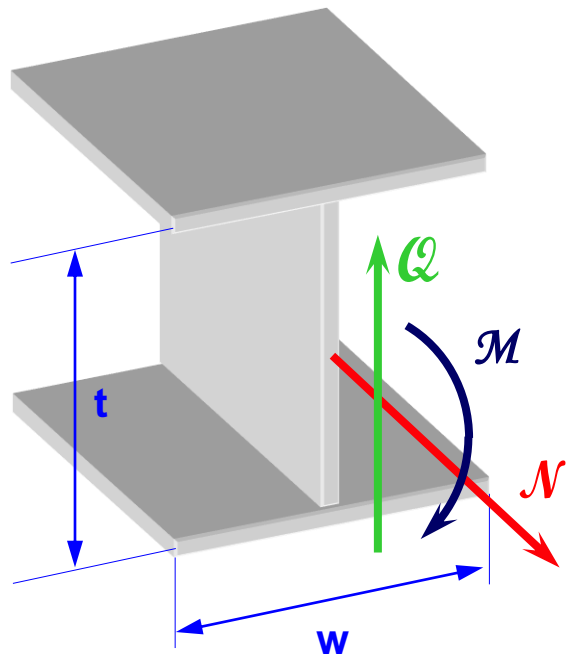


Chapter 4. Composite beams and plates

Composite beams

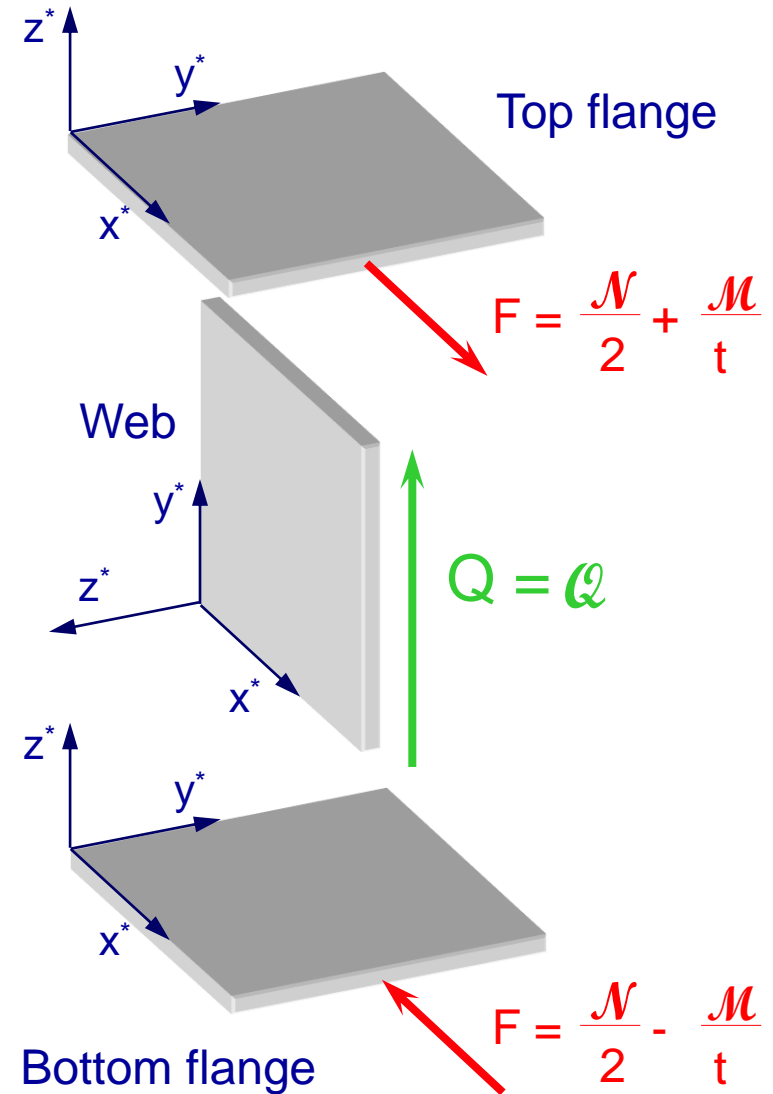
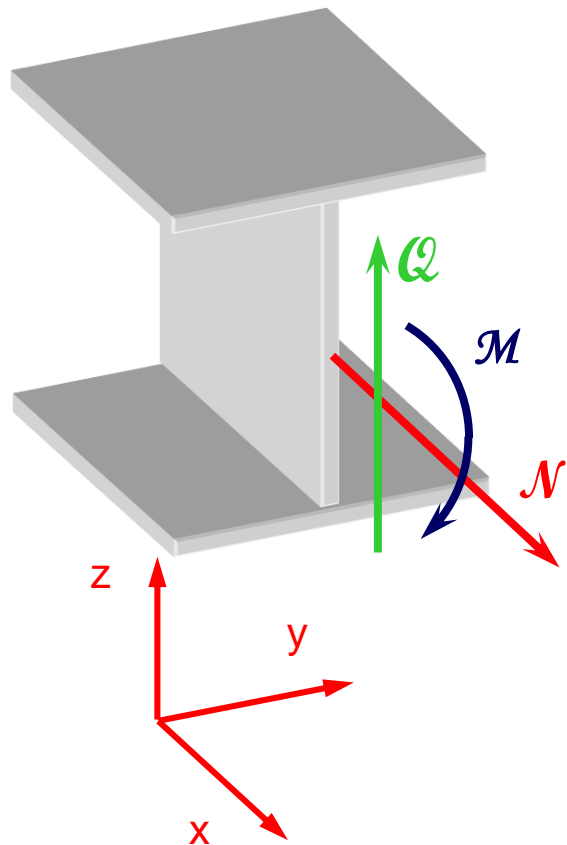
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Bending of beams



Stresses in thin-walled beams

Bending of beams



Stresses in thin-walled beams

Bending of beams

Top flange

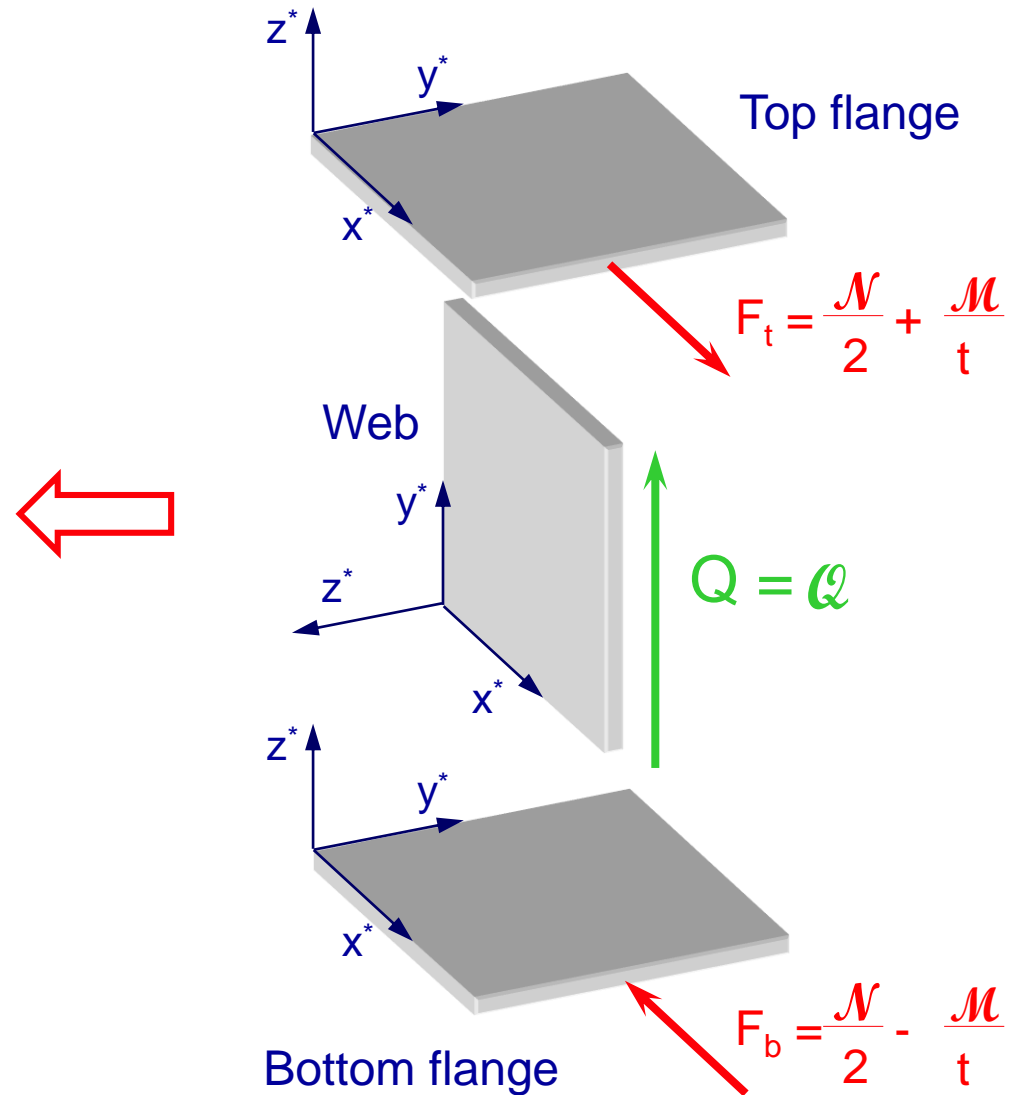
$$\{N\} = \begin{Bmatrix} N_x \\ N_y \\ N_{xy} \end{Bmatrix} = \begin{Bmatrix} F_t/w \\ 0 \\ 0 \end{Bmatrix}$$

Web

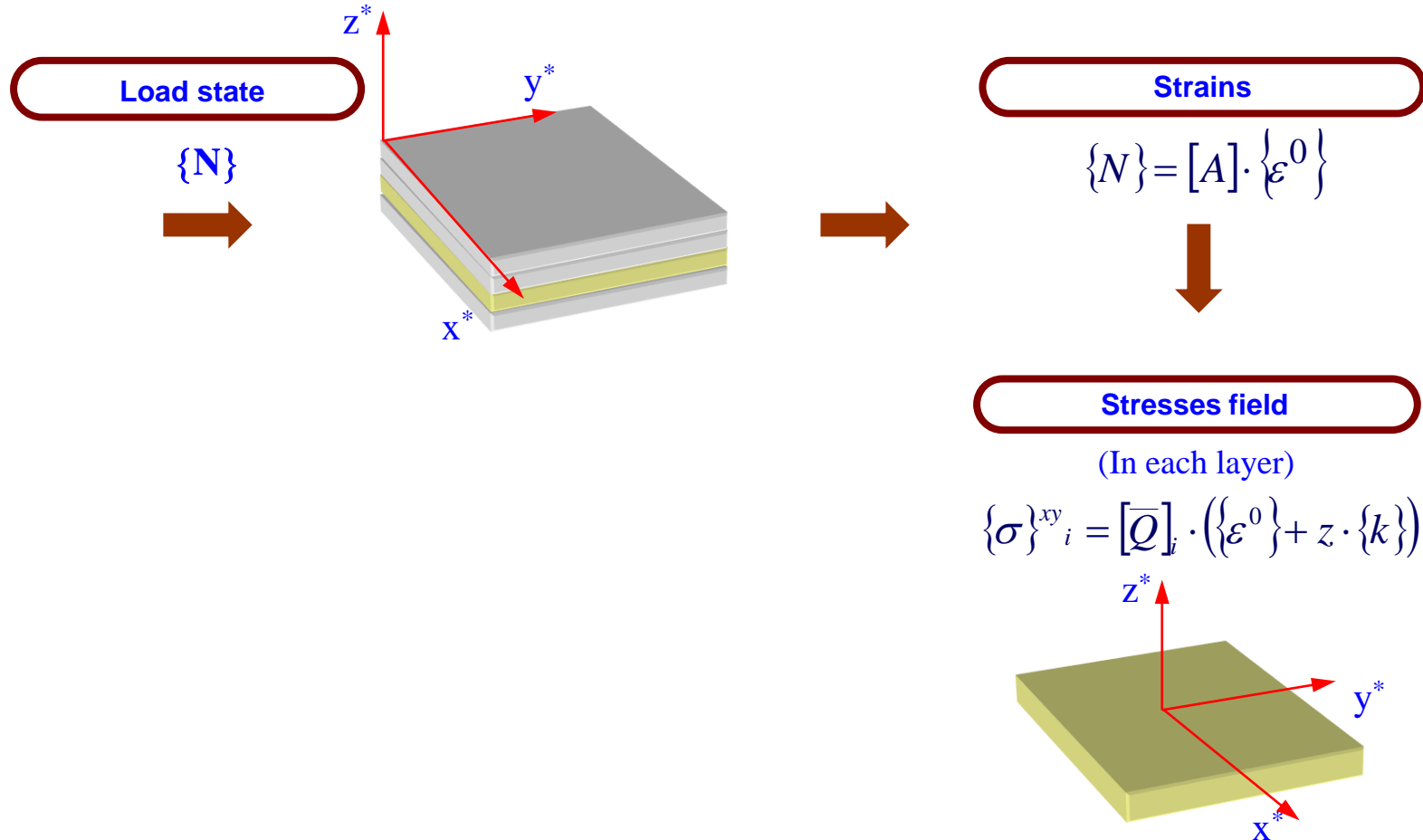
$$\{N\} = \begin{Bmatrix} N_x \\ N_y \\ N_{xy} \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ Q/t \end{Bmatrix}$$

Bottom flange

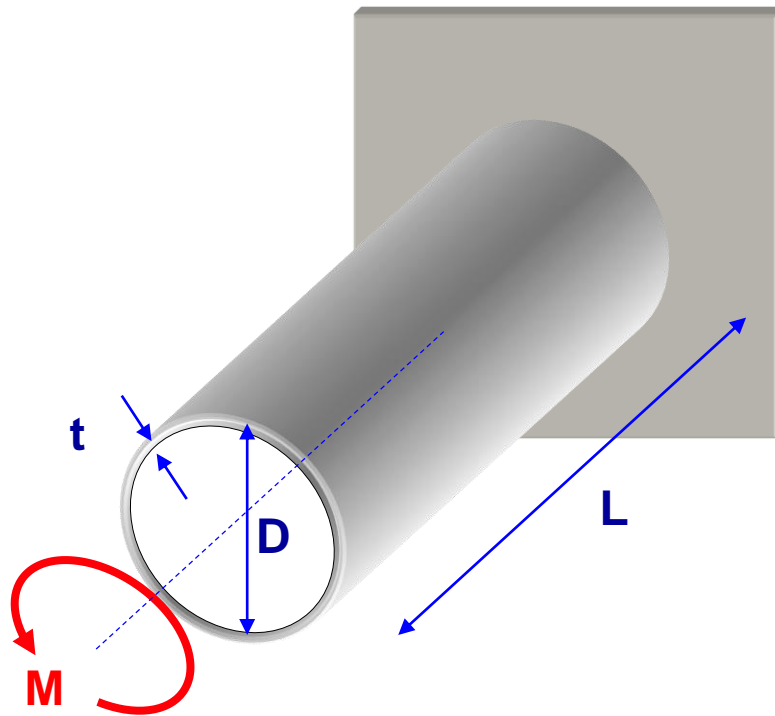
$$\{N\} = \begin{Bmatrix} N_x \\ N_y \\ N_{xy} \end{Bmatrix} = \begin{Bmatrix} F_b/w \\ 0 \\ 0 \end{Bmatrix}$$



Bending of beams



Torsion of beams



Hypothesis:

$$H \ll D$$

Torsion of beams

Internal forces

Theory for isotropic materials:

$$\tau_{xy}^0 = \frac{2 \cdot M}{\pi \cdot D^2 \cdot t}$$



$$\begin{Bmatrix} N_x \\ N_y \\ N_{xy} \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ \frac{2 \cdot M}{D \cdot t} \end{Bmatrix}$$

Torsion of beams

Strains

$$\begin{Bmatrix} \varepsilon_x^o \\ \varepsilon_y^o \\ \gamma_{xy}^o \end{Bmatrix} = \begin{bmatrix} a_{11} & a_{12} & a_{1s} \\ a_{12} & a_{22} & a_{2s} \\ a_{1s} & a_{2s} & a_{ss} \end{bmatrix} \begin{Bmatrix} 0 \\ 0 \\ \frac{2 \cdot M}{D \cdot t} \end{Bmatrix}$$

En particular

$$\gamma_{xy}^o = a_{ss} \cdot \frac{2 \cdot M}{D \cdot t}$$

Torsion of beams

Strains

$$\gamma_{xy}^o = a_{ss} \cdot \frac{2 \cdot M}{D \cdot t}$$

$$\gamma_{xy}^o = \frac{2 \cdot M}{\pi \cdot D^2 \cdot t \cdot G_{xy}^o} \quad G_{xy}^o = \frac{\pi \cdot D}{a_{ss}}$$

Rotations:

$$\theta_{\max} = \frac{M \cdot L}{G_{xy}^o \cdot I_o} \quad I_o = \frac{\pi \cdot D^3 \cdot t}{4}$$

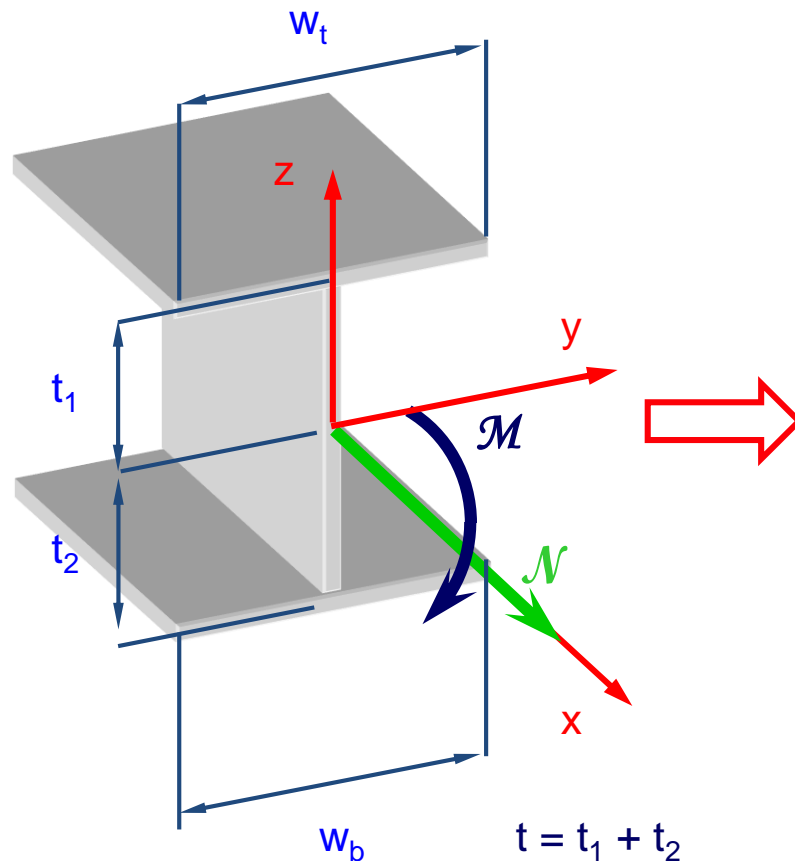


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Axial force and bending moment



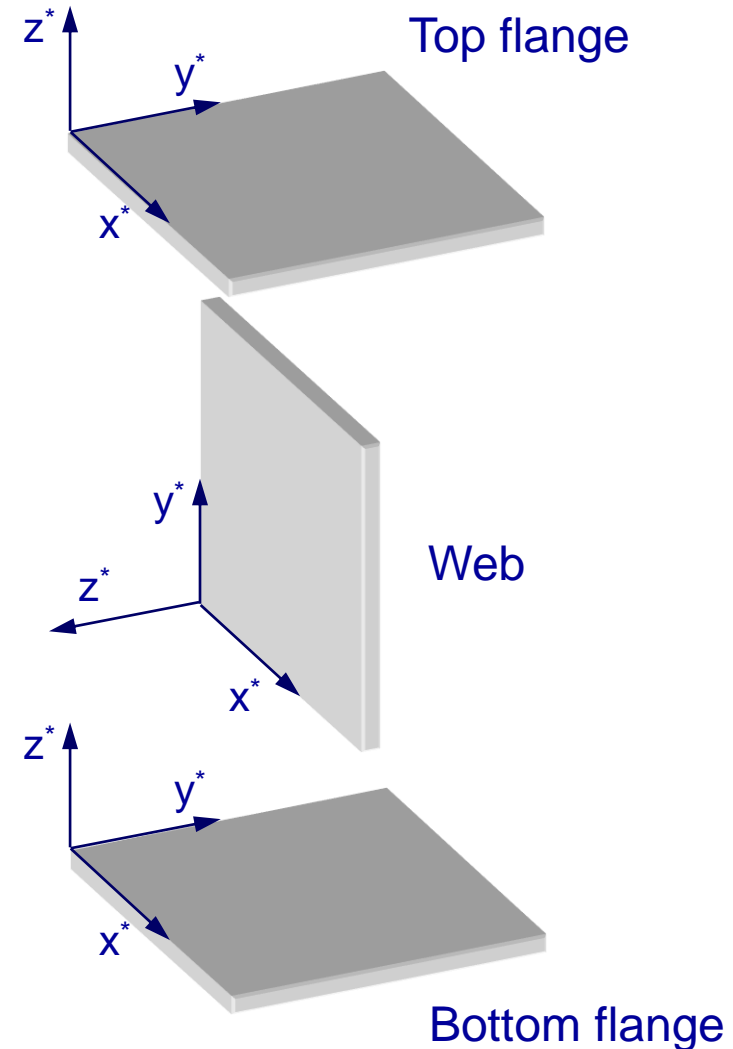
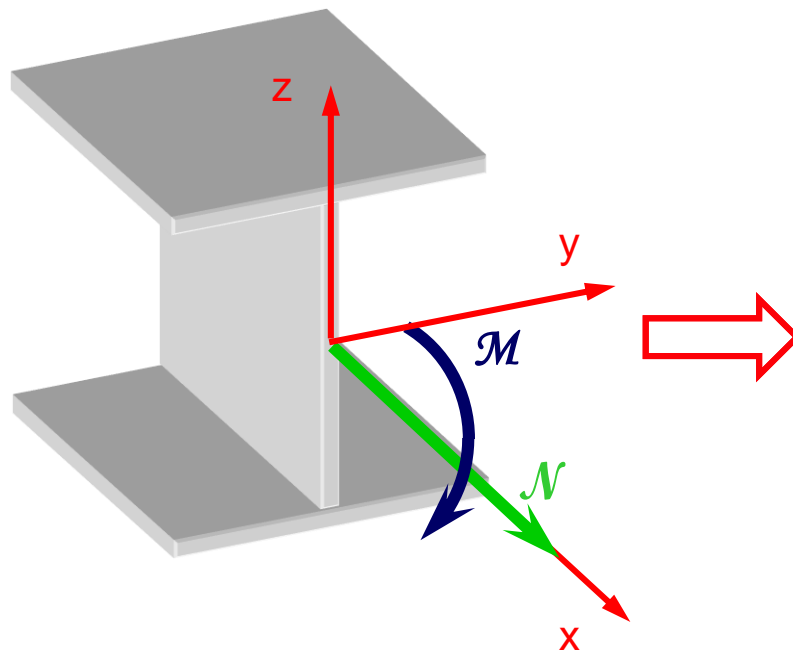
Hypotheses

- The bending moment can be found through the integral of the normal stress along the z direction.
- Longitudinal strain in x direction is directly related to the distance from neutral axis.

$$\varepsilon_x(z) = z \cdot k_x$$

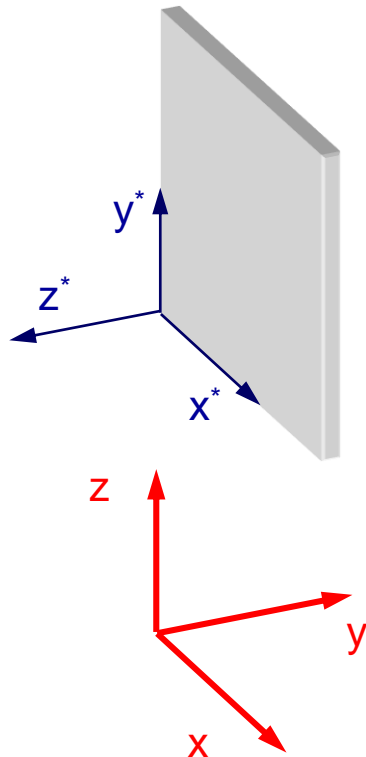
Displacements in thin-walled beams

Axial force and bending moment

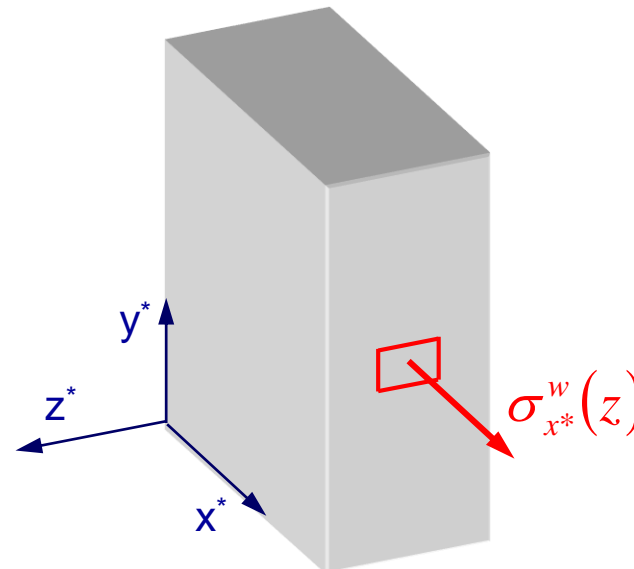


Axial force and bending moment

Web

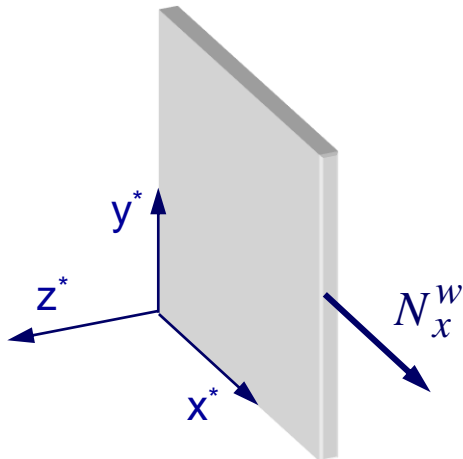


- Bending stiffness matrix $[D]$ is not necessary.
- Web properties do not depend on z .
- Stresses on the web depend on z .



Axial force and bending moment

Web



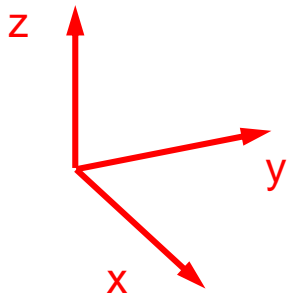
$$N_{x^*}^w = \int \sigma_{x^*}^w \cdot dz^*$$

$$\begin{Bmatrix} \varepsilon_{x^*}^{o^w} \\ \varepsilon_{y^*}^{o^w} \\ \gamma_{xy^*}^{o^w} \end{Bmatrix} = \begin{bmatrix} a_{11^*}^w & a_{11^*}^w & a_{11^*}^w \\ a_{11^*}^w & a_{11^*}^w & a_{11^*}^w \\ a_{11^*}^w & a_{11^*}^w & a_{11^*}^w \end{bmatrix} \cdot \begin{Bmatrix} N_{x^*}^w \\ 0 \\ 0 \end{Bmatrix}$$

$$\varepsilon_{x^*}^{o^w} = a_{11^*}^w \cdot N_{x^*}^w$$



$$\varepsilon_x^{o^w}(z) = a_{11^*}^w \cdot N_x^w(z)$$

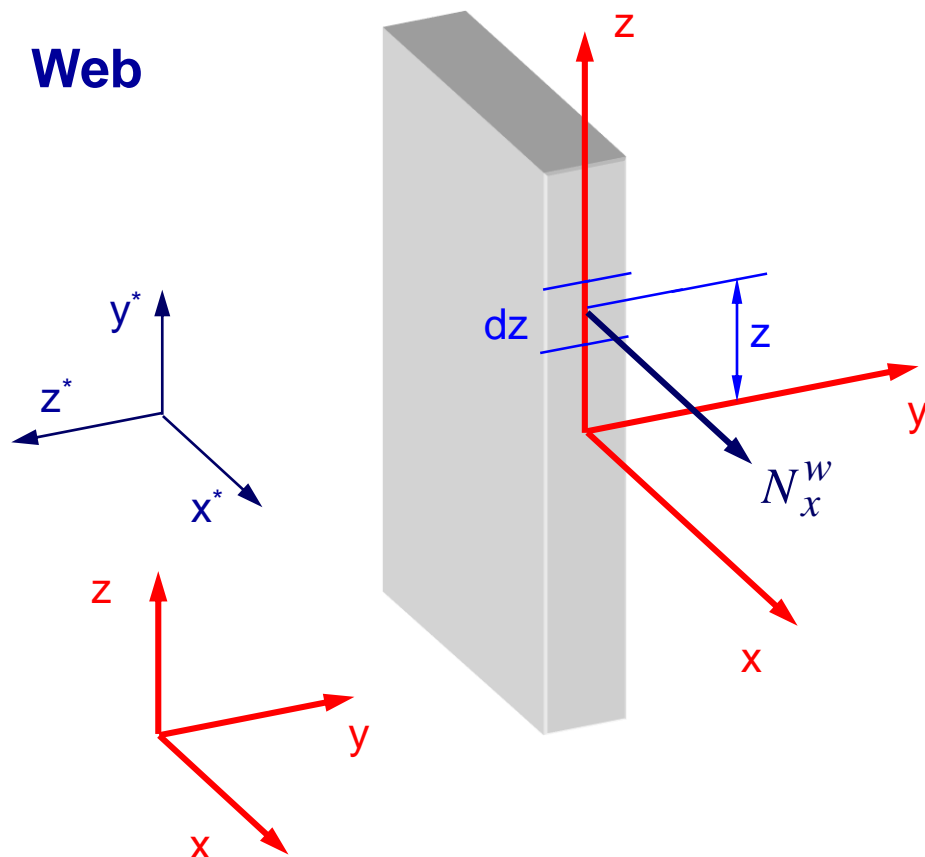


! Axes x^* and x
match up !

! Results depend
on global
coordinate z !

Displacements in thin-walled beams

Axial force and bending moment



$$\varepsilon_x^{0w}(z) = z \cdot k_x$$

**Axial force per
unit length,
coordinate y^*
(coordinate z)**

$$N_x^w = \frac{1}{a_{11}^w} \cdot z \cdot k_x$$

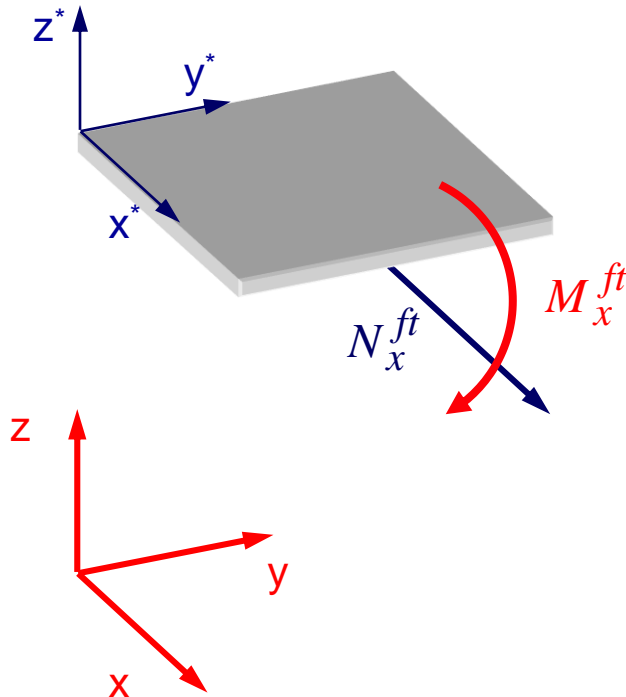
This axial force contributes to
the section bending moment

$$\mathcal{M}_w = \int_{-t/2}^{t/2} N_x^w \cdot z \cdot dz$$

$$\mathcal{M}_w = \frac{t^3}{12 \cdot a_{11}^w} \cdot k_x$$

Axial force and bending moment

Top flange



$$\begin{Bmatrix} \varepsilon_{x^*}^{o\,ft} \\ \varepsilon_{y^*}^{o\,ft} \\ \gamma_{xy^*}^{o\,ft} \end{Bmatrix} = \begin{bmatrix} a_{11^*}^{ft} & a_{11^*}^{ft} & a_{11^*}^{ft} \\ a_{11^*}^{ft} & a_{11^*}^{ft} & a_{11^*}^{ft} \\ a_{11^*}^{ft} & a_{11^*}^{ft} & a_{11^*}^{ft} \end{bmatrix} \cdot \begin{Bmatrix} N_{x^*}^{ft} \\ 0 \\ 0 \end{Bmatrix}$$

$$\begin{Bmatrix} k_{x^*} \\ k_{y^*} \\ k_{xy^*} \end{Bmatrix} = \begin{bmatrix} d_{11^*}^{ft} & d_{11^*}^{ft} & d_{11^*}^{ft} \\ d_{11^*}^{ft} & d_{11^*}^{ft} & d_{11^*}^{ft} \\ d_{11^*}^{ft} & d_{11^*}^{ft} & d_{11^*}^{ft} \end{bmatrix} \cdot \begin{Bmatrix} M_{x^*}^{ft} \\ 0 \\ 0 \end{Bmatrix}$$

$$\varepsilon_{x^*}^{o\,ft} = a_{11^*}^{ft} \cdot N_{x^*}^{ft}$$

$$k_{x^*} = d_{11^*}^{ft} \cdot M_{x^*}^{ft}$$



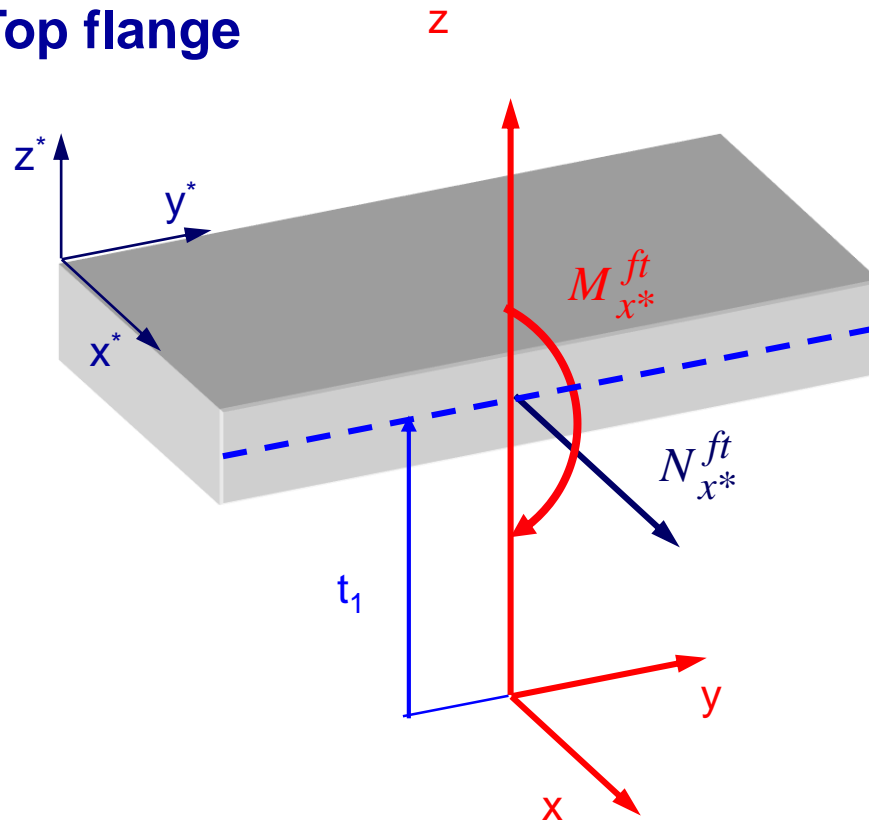
$$\varepsilon_x^{o\,ft} = a_{11^*}^{ft} \cdot N_x^{ft}$$

$$k_x = d_{11^*}^{ft} \cdot M_x^{ft}$$

Displacements in thin-walled beams

Axial force and bending moment

Top flange



Axial force and bending moment in the flange contribute to the section bending moment

$$\mathcal{M}_{ft} = w_t \cdot (N_x^{ft} \cdot t_1 + M_x^{ft})$$

$$\mathcal{M}_{ft} = w_t \cdot \left(\frac{t_1}{a_{11}^{ft}} \cdot \varepsilon_x^{o_{ft}} + \frac{1}{d_{11}^{ft}} \cdot k_x \right)$$

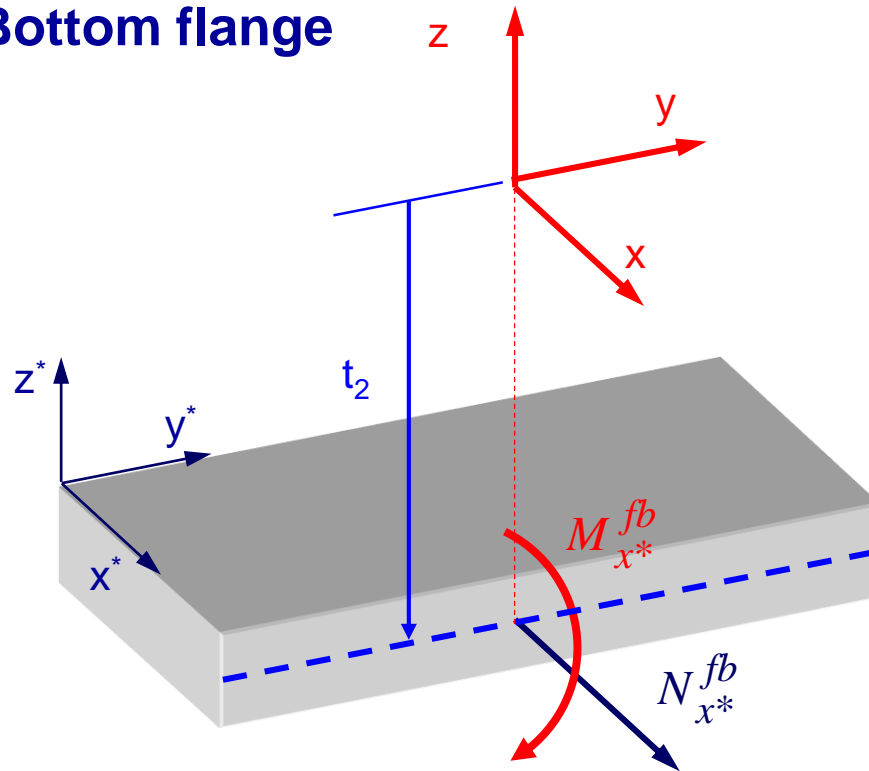
Strains at flange mid-plane

$$\varepsilon_x^{o_{ft}} = t_1 \cdot k_x$$

$$\mathcal{M}_{ft} = w_t \cdot \left(\frac{t_1^2}{a_{11}^{ft}} + \frac{1}{d_{11}^{ft}} \right) \cdot k_x$$

Axial force and bending moment

Bottom flange



By analogy:

$$\mathcal{M}_{fb} = w_b \cdot \left(\frac{t_2^2}{a_{11}^{fb}} + \frac{1}{d_{11}^{fb}} \right) \cdot k_x$$

Displacements in thin-walled beams

Axial force and bending moment

The section bending moment is the sum of the bending moments supported by web and flanges:

$$\mathcal{M} = \left[w_t \cdot \left(\frac{t_1^2}{a_{11}^{ft}} + \frac{1}{d_{11}^{ft}} \right) + w_b \cdot \left(\frac{t_2^2}{a_{11}^{fb}} + \frac{1}{d_{11}^{fb}} \right) + \frac{(t_1 + t_2)^3}{12 \cdot a_{11}^w} \right] \cdot k_x$$

In isotropic material beams:

$$\mathcal{M} = EI \cdot k_x$$

$$\langle EI \rangle = w_t \cdot \left(\frac{t_1^2}{a_{11}^{ft}} + \frac{1}{d_{11}^{ft}} \right) + w_b \cdot \left(\frac{t_2^2}{a_{11}^{fb}} + \frac{1}{d_{11}^{fb}} \right) + \frac{(t_1 + t_2)^3}{12 \cdot a_{11}^w}$$



Chapter 4. Composite beams and plates

Composite beams

- 1. Introduction**
- 2. Analysis of beams**
- 3. Stresses in solid beams**
- 4. Displacements in solid beams**
- 5. Stresses in thin-walled beams**
- 6. Displacements in thin-walled beams**
- 7. References**



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