



Universidad
Carlos III de Madrid
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**Department of Continuum Mechanics and Structural
Engineering**

Aerospace Structures

Chapter 4. Plates and Shells

Circular plates



Chapter 4. Plates and Shells

Circular plates

1. Equilibrium equations in polar coordinate system
2. Bending equations in polar coordinate system
3. Boundary conditions in polar coordinate system
4. Examples
5. References



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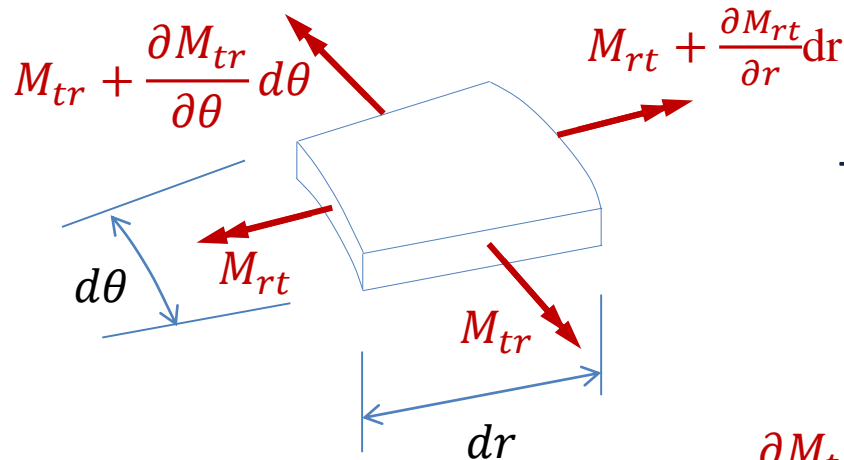
- Bending and torsion moments:

Defined per unit length

$$M_r = \int_{-\frac{h}{2}}^{+\frac{h}{2}} \sigma_{rr} \cdot z \cdot dz$$

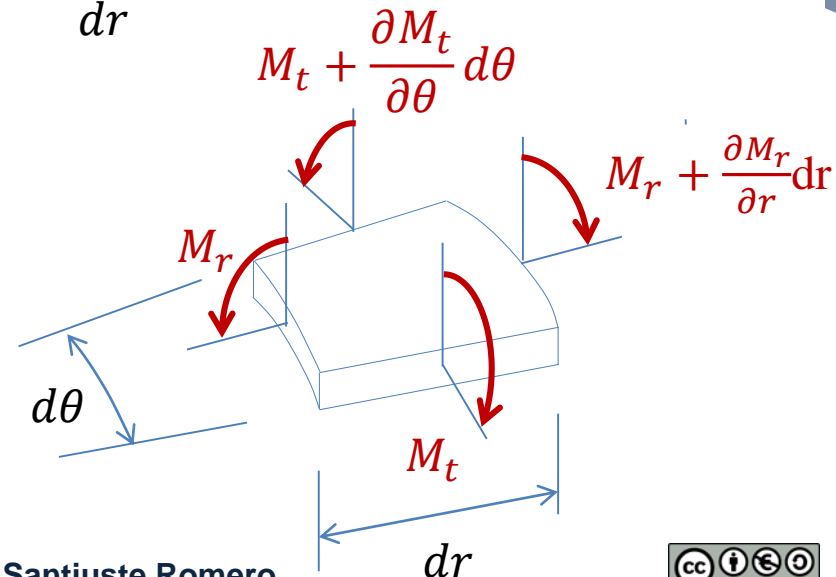
$$M_t = \int_{-\frac{h}{2}}^{+\frac{h}{2}} \sigma_{\theta\theta} \cdot z \cdot dz$$

$$M_{rt} = \int_{-\frac{h}{2}}^{+\frac{h}{2}} \sigma_{r\theta} \cdot z \cdot dz$$



Torsion moment

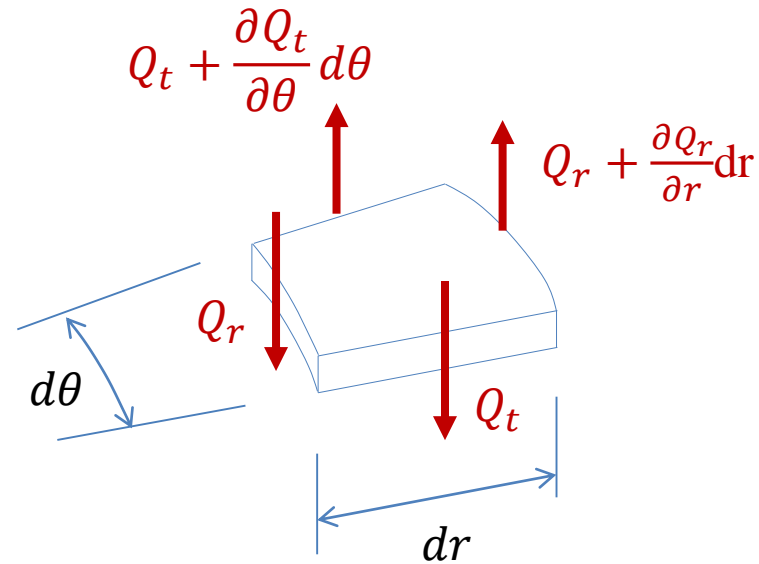
Bending moments



- Shear force

$$Q_r = \int_{-\frac{h}{2}}^{+\frac{h}{2}} \tau_{rz} \cdot z dz$$

$$Q_t = \int_{-\frac{h}{2}}^{+\frac{h}{2}} \tau_{\theta z} \cdot z dz$$



Defined per unit length

Internal forces can be expressed in terms of vertical displacement (same procedure than in Cartesian coordinate system)

$$M_r = -D \cdot \left(\frac{\partial^2 w}{\partial r^2} + \nu \cdot \left(\frac{1}{r} \frac{\partial w}{\partial r} + \frac{1}{r^2} \frac{\partial^2 w}{\partial \theta^2} \right) \right)$$

$$M_t = -D \cdot \left(\frac{1}{r} \frac{\partial w}{\partial r} + \frac{1}{r^2} \frac{\partial^2 w}{\partial \theta^2} + \nu \frac{\partial^2 w}{\partial r^2} \right)$$

$$M_{rt} = -D \cdot (1 - \nu) \left(\frac{1}{r} \frac{\partial^2 w}{\partial r \partial \theta} - \frac{1}{r^2} \frac{\partial w}{\partial \theta} \right)$$

$$Q_r = -D \cdot \frac{\partial}{\partial r} (\nabla^2 w)$$

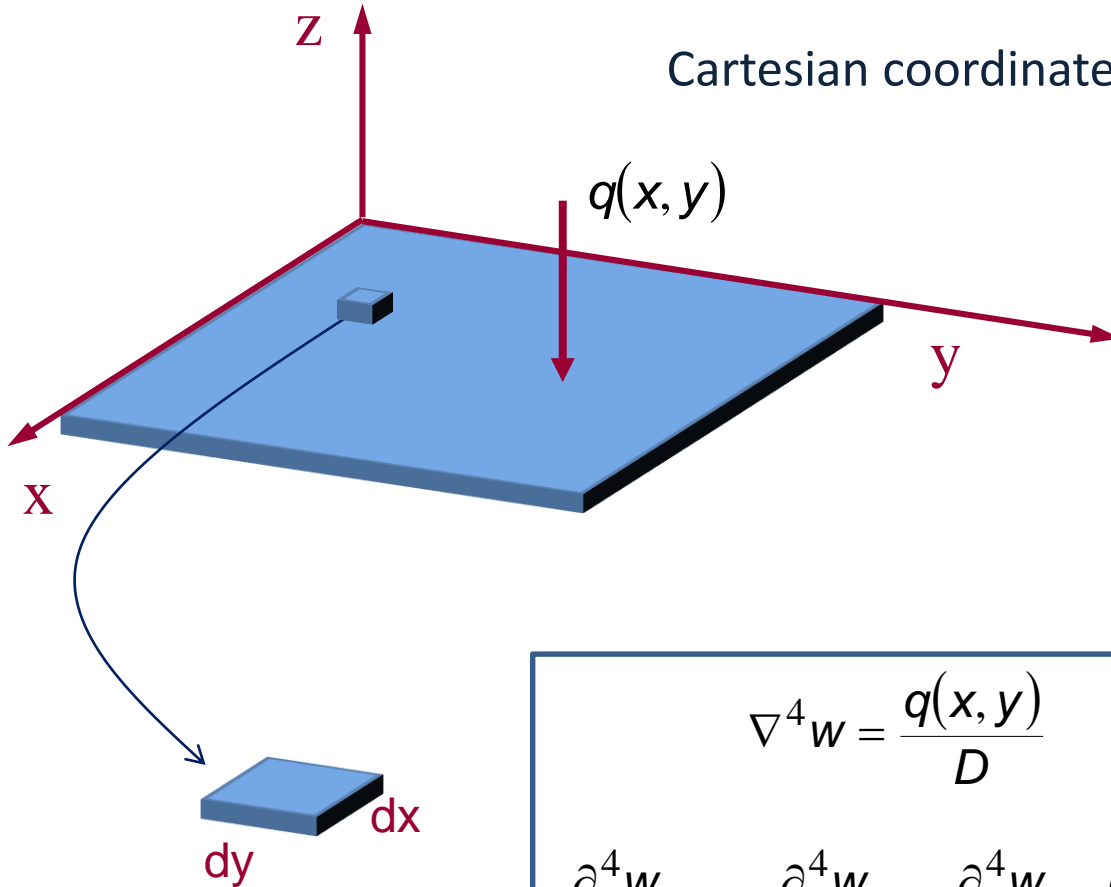
$$Q_t = -\frac{D}{r} \cdot \frac{\partial}{\partial \theta} (\nabla^2 w)$$



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Cartesian coordinate system:

Del operator

$$\nabla = \frac{\partial}{\partial x} \vec{i} + \frac{\partial}{\partial y} \vec{j} + \frac{\partial}{\partial z} \vec{k}$$

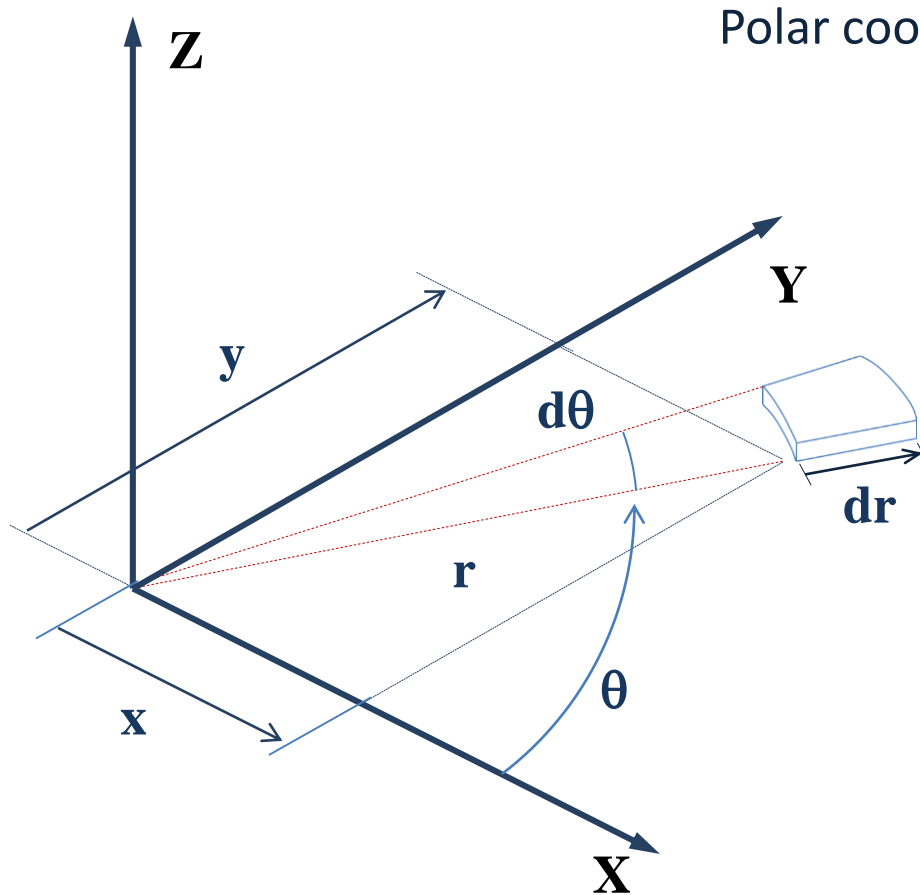
Laplace operator (in-plane)

$$\nabla^2 = \Delta = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$$

$$\nabla^4 w = \frac{q(x, y)}{D}$$

$$\frac{\partial^4 w}{\partial x^4} + 2 \frac{\partial^4 w}{\partial x^2 \partial y^2} + \frac{\partial^4 w}{\partial y^4} = \frac{q(x, y)}{D}$$

$$D = \frac{E \cdot h^3}{12 \cdot (1 - \nu^2)}$$

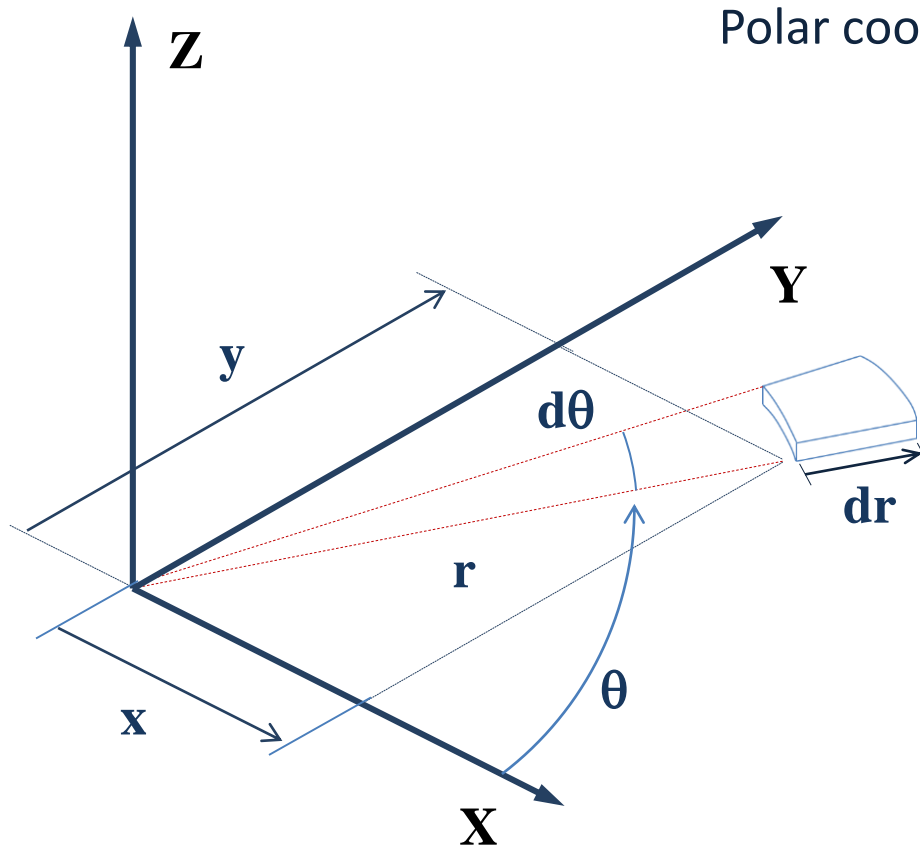


Del operator

$$\nabla = \frac{\partial}{\partial r} \bar{u}_r + \frac{1}{r} \frac{\partial}{\partial \theta} \bar{u}_\theta + \frac{\partial}{\partial z} \bar{u}_z$$

Laplace operator (in-plane)

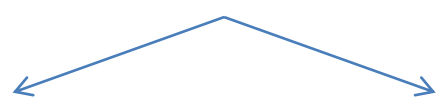
$$\nabla^2 = \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2}$$



$$\nabla^4 w = \frac{q}{D}$$

$$\nabla^4 w(r, \theta) = \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} \right) \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} \right) w = \frac{q(r, \theta)}{D}$$

Solution of the differential equation

$$W = W_h + W_p$$

$$\nabla^4 W_h = 0 \qquad \nabla^4 W_p = \frac{q}{D}$$

The solution of the homogeneous equation can be expressed as the next series:

$$w_h(r, \theta) = \sum_{m=0}^{\infty} R_m(r) \cdot \cos(m \cdot \theta) + \sum_{m=1}^{\infty} S_m(r) \cdot \text{sen}(m \cdot \theta)$$

$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} \right) \left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} \right) R_m = 0$$

$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} \right) \left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{m^2}{r^2} \right) S_m = 0$$

The solutions to these equations are:

$$R_o = A_o + B_o \cdot \ln r + C_o \cdot r^2 + D_o \cdot r^2 \cdot \ln r$$

$$R_1, S_1 = A_1 \cdot r + B_1 \cdot \frac{1}{r} + C_1 \cdot r^3 + D_1 \cdot r \cdot \ln r$$

$$R_m, S_m = A_m \cdot r^m + B_m \cdot \frac{1}{r^m} + C_m \cdot r^{m+2} + D_m \cdot \frac{1}{r^{m-2}}$$

The coefficient of these equations can be found imposing boundary conditions



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Clamped in $r = a$

$$w|_{r=a} = 0$$

$$\frac{\partial w}{\partial r}|_{r=a} = 0$$

Simply supported in $r = a$

$$w|_{r=a} = 0$$

$$M_r|_{r=a} = 0$$

Free end in $r = a$

$$\left(Q_r + \frac{1}{r} \cdot \frac{\partial M_{tr}}{\partial \theta} \right)|_{r=a} = 0$$

$$M_r|_{r=a} = 0$$

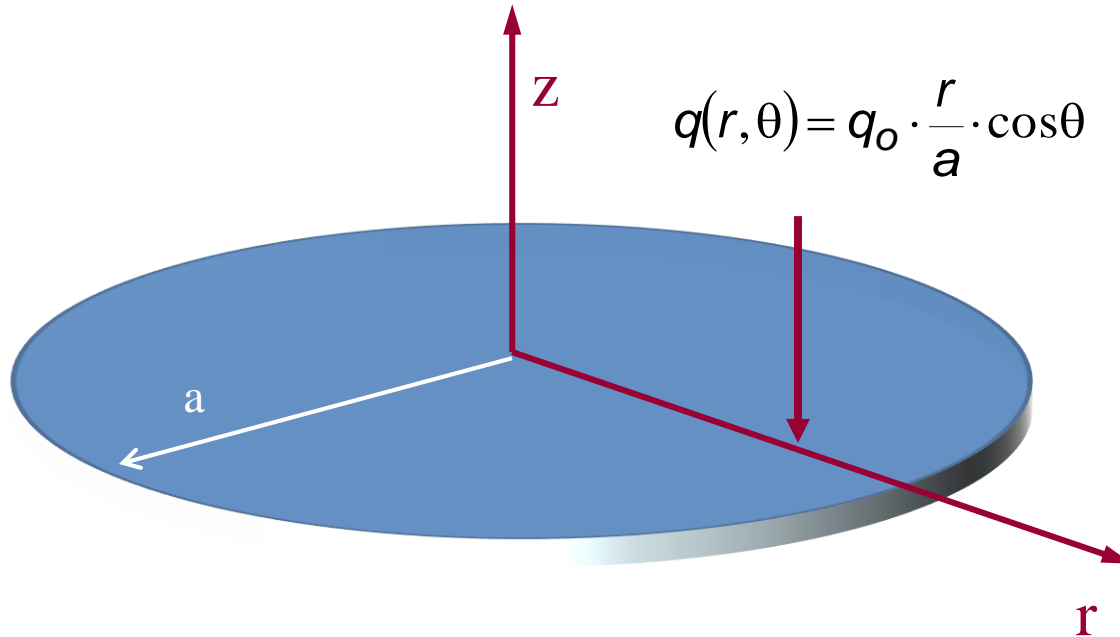


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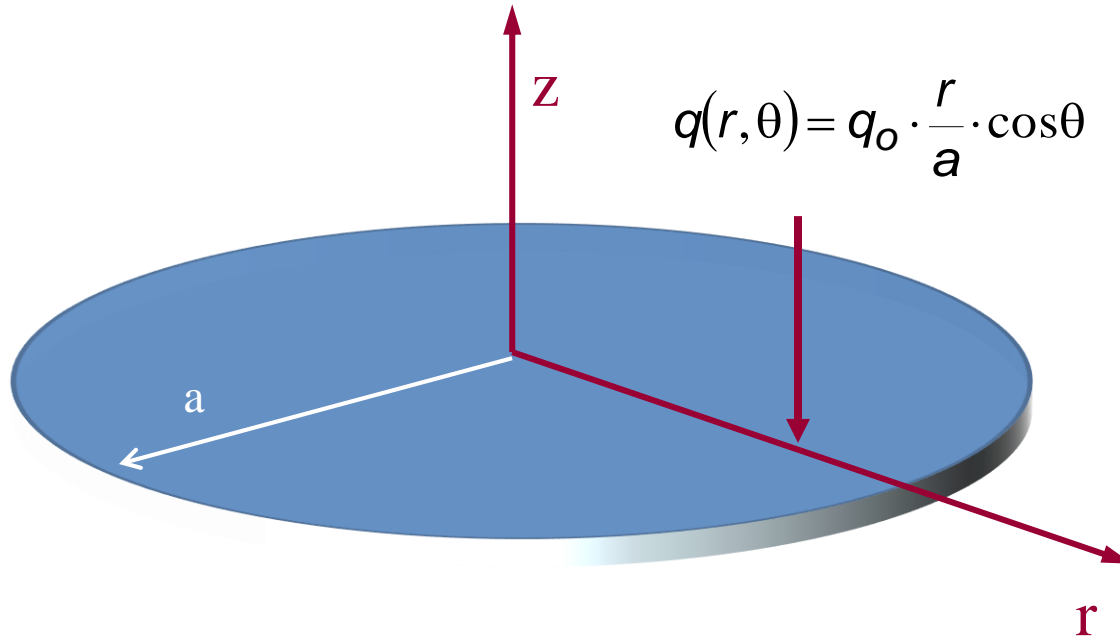
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Example: Load directly related to radius



$$\left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} \right) \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} \right) w = \frac{q_0 \cdot r \cdot \cos\theta}{a \cdot D}$$

Example: Load directly related to radius



Particular solution

$$\nabla^4 w_p = \frac{q_0 \cdot r \cdot \cos\theta}{a \cdot D}$$



There are fourth derivatives in the equation, thus this is a possible solution:

Where A is a constant

$$w_p(r, \theta) = A \cdot r^5 \cdot \cos\theta$$

Example: Load directly related to radius

$$w_p(r, \theta) = A \cdot r^5 \cdot \cos \theta$$

$$\left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} \right) \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} \right) w = \frac{q_0 \cdot r \cdot \cos \theta}{a \cdot D}$$

$$w_p(r, \theta) = 5.2 \cdot 10^{-3} \cdot \frac{q_0 \cdot r^5 \cdot \cos \theta}{a \cdot D}$$

Particular solution

Example: Load directly related to radius

Homogenous solution

$$w_h(r, \theta) = \sum_{m=0}^{\infty} R_m(r) \cdot \cos(m \cdot \theta) + \sum_{m=1}^{\infty} S_m(r) \cdot \text{sen}(m \cdot \theta)$$

Since load only has terms in cosine, the solution is:

$$w_h(r, \theta) = R_1(r) \cdot \cos \theta$$
$$R_1(r) = A_1 \cdot r + B_1 \cdot \frac{1}{r} + C_1 \cdot r^3 + D_1 \cdot r \cdot \ln r$$

Example: Load directly related to radius

Final solution:

$$w(r, \theta) = \frac{q_0 \cdot a^4}{192 \cdot D} \left(\rho^5 + A^* \cdot \rho + B^* \cdot \frac{1}{\rho} + C^* \cdot \rho^3 + D^* \cdot \rho \cdot \ln \rho \right) \cdot \cos \theta \quad \rho = \frac{r}{a}$$

Constant are calculated imposing boundary conditions

- Vertical displacement is null in the midpoint:

$$B^* = 0$$

Example: Load directly related to radius

$$M_r = -D \cdot \left(\frac{\partial^2 w}{\partial r^2} + \nu \cdot \left(\frac{1}{r} \frac{\partial w}{\partial r} + \frac{1}{r^2} \frac{\partial^2 w}{\partial \theta^2} \right) \right)$$

$$M_t = -D \cdot \left(\frac{1}{r} \frac{\partial w}{\partial r} + \frac{1}{r^2} \frac{\partial^2 w}{\partial \theta^2} + \nu \frac{\partial^2 w}{\partial r^2} \right)$$

$$M_{rt} = -D \cdot (1 - \nu) \left(\frac{1}{r} \frac{\partial^2 w}{\partial r \partial \theta} - \frac{1}{r^2} \frac{\partial w}{\partial \theta} \right)$$



$$M_r = -D \cdot \frac{q_0 \cdot a^2}{192} \left[(20 + 4 \cdot \nu) \cdot \rho^3 + (6 + 2 \cdot \nu) \cdot C^* \cdot \rho + \frac{D^* \cdot (1 + \nu)}{\rho} \right] \cdot \cos \theta$$

$D^* = 0$ Otherwise bending moment is infinite in the midpoint

Example: Load directly related to radius

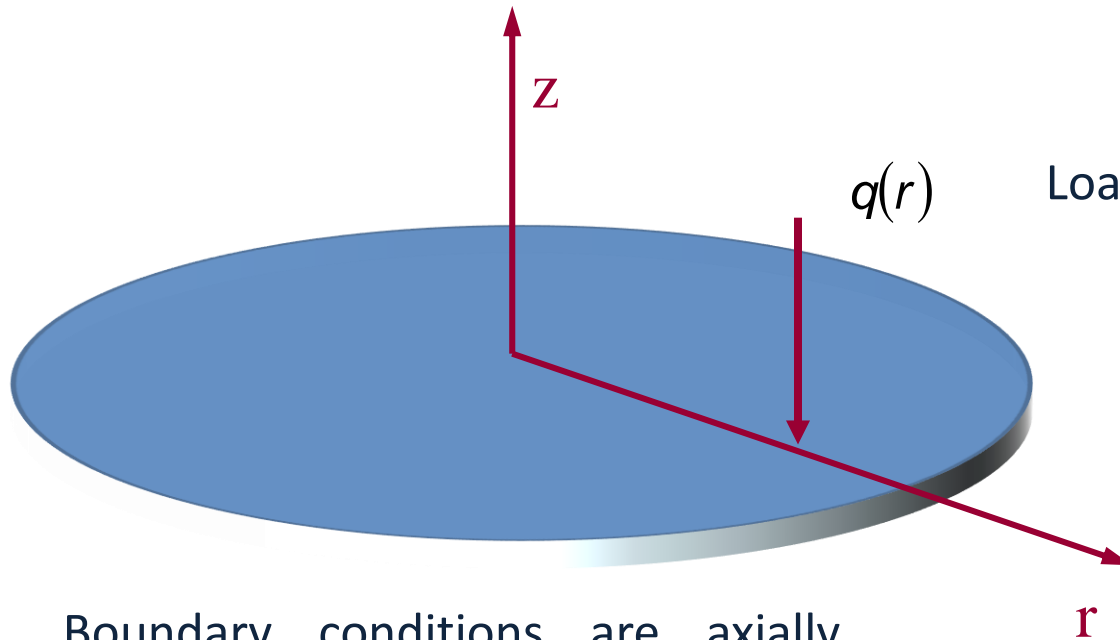
$$w(r, \theta) = \frac{q_0 \cdot a^4}{192 \cdot D} \left(\frac{r^5}{a^5} + A^* \cdot \frac{r}{a} + C^* \cdot \frac{r^3}{a^3} \right) \cdot \cos \theta$$

The other constants are found imposing boundary conditions at the external edge $r=a$

Example, simply supported plate:

$$\begin{array}{l} w|_{r=a} = 0 \\ M_r|_{r=a} = 0 \end{array} \quad \Rightarrow \quad \begin{array}{l} A^* = \frac{7+\nu}{3+\nu} \\ C^* = -2 \frac{5+\nu}{3+\nu} \end{array}$$

Example: Axial symmetry



Loads are angle independent

Boundary conditions are axially symmetrical

$$\nabla^4 w = \frac{q}{D}$$

$$\nabla^4 w(r) = \left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{\partial}{\partial r} \right) \left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{\partial}{\partial r} \right) w = \frac{q(r)}{D}$$

Example: Axial symmetry

Internal forces equations considering angle independency

$$M_r = -D \cdot \left(\frac{d^2 w}{dr^2} + \frac{\nu}{r} \frac{dw}{dr} \right)$$

$$M_t = -D \cdot \left(\frac{1}{r} \frac{dw}{dr} + \nu \frac{d^2 w}{dr^2} \right)$$

$$M_{rt} = 0$$

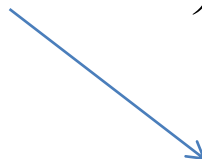
$$Q_r = -D \cdot \frac{d}{dr} \left(\frac{1}{r} \frac{d}{dr} \left(r \frac{dw}{dr} \right) \right)$$

$$Q_t = 0$$

Example: Axial symmetry

Differential equation can be integrated when load is uniform:

$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{\partial}{\partial r} \right) \left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{\partial}{\partial r} \right) = \frac{q_o}{D}$$


$$w_h(r) = A_o + B_o \cdot \ln r + C_o \cdot r^2 + D_o \cdot r^2 \cdot \ln r$$

$$w_p(r) = \frac{q_o \cdot r^4}{64D}$$

Coefficients are found imposing boundary conditions



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